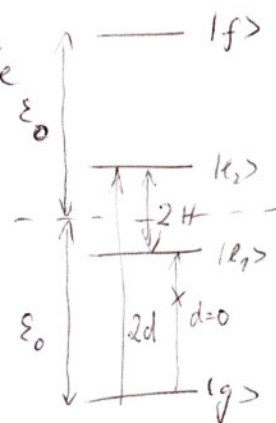


Let us investigate simple model systems with simple models for the energy gap correlation function

Relaxation in H-aggregates - studied by pump-probe

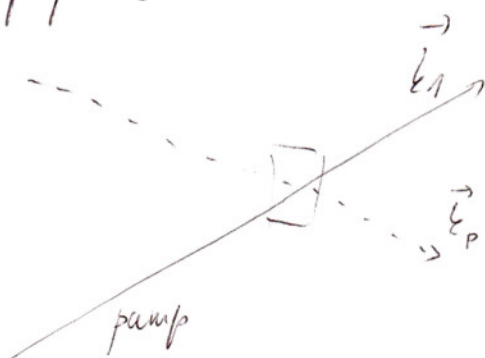


First three excited states would look like



Let us measure the life time of the state $|l_2\rangle$

by pump probe



Absorption without pump

$$H(\omega) \approx \int_0^{\infty} dt e^{-g(t) + i(\omega - \omega_{12})t}$$

Absorption of the probe after pulse can be calculated from $P(\omega)$ that is a result of interaction with $-\vec{\epsilon}_1 + \vec{\epsilon}_2 + \vec{\epsilon}_p$ light

$\vec{\epsilon}_1, \vec{\epsilon}_2$ 2 pulses at the same time
 $t_1 = 0$

$$P(t, T, 0) \approx -i [\bar{R}_2(t, T, 0) + \bar{R}_3(t, T, 0)] e^{-i\omega t} E_p^3$$

relaxation
↓

$$\bar{R}_2(t, T, 0) \approx e^{-g^*(t) + g^*(T) - g(T)} - [g^*(T+t) - g(T+t)] e^{-KT}$$

$g \rightarrow g - i\Gamma$ $- (g + i\Gamma)$

$$\bar{R}_3(t, T, 0) \approx e^{-g(t)}$$

$$\bar{R}_1(t, T, 0) = i |d|^4 \frac{1}{\epsilon} e^{-g(t)} \frac{1}{\epsilon} e^{-i\omega_0 t}$$

$$\bar{R}_2(t, T, 0) = i |d|^4 \frac{1}{\epsilon} e^{-g^*(t)} \frac{1}{\epsilon} \text{Im} \{g(T) - g(T+t)\} \frac{1}{\epsilon} e^{-kt} \frac{1}{\epsilon} e^{-i\omega_0 t}$$

$$g(t) = \Gamma t$$



$$\bar{R}_3(t, T, 0) = |d|^4 \frac{1}{\epsilon} e^{-\Gamma t}$$

$$\bar{R}_4(t, T, 0) = |d|^4 \frac{1}{\epsilon} e^{-\Gamma t} \frac{1}{\epsilon} e^{-kT}$$

$$P(t, T, 0) = -i |d|^4 \frac{1}{\epsilon} e^{-\Gamma t} (1 + e^{-kT}) \frac{1}{\epsilon} e^{-i\omega_0 t}$$

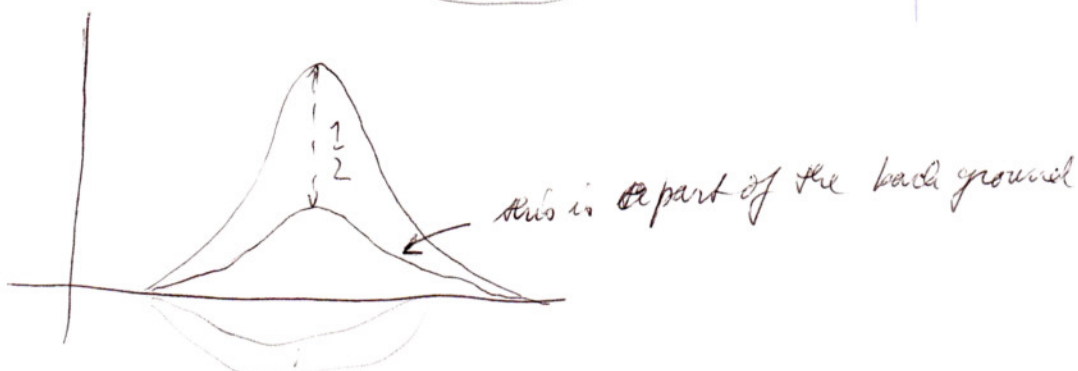
$$P(\omega, T, 0) = -i |d|^4 (1 + e^{-kT}) \int_{-\infty}^{\infty} dt \frac{1}{\epsilon} e^{-\Gamma t + i(\omega - \omega_0)t} \quad \text{②} \quad E_{\text{pump}} E_0$$

$$= -i |d|^4 (1 + e^{-kT}) \left[\frac{-1}{\Gamma - i(\omega - \omega_0)} \right]_0^{\infty} E_{\text{pump}}^2 E_0 = i |d|^4 \frac{1 (1 + e^{-kT})}{\Gamma - i(\omega - \omega_0)}$$

$$= i |d|^4 \frac{(1 + e^{-kT})}{\Gamma^2 - (\omega - \omega_0)^2} (\Gamma + i(\omega - \omega_0)) \quad \times E_{\text{pump}}^2 E_0$$

$$E_p(t) = E_0 \delta(t) e^{-i\omega t} \Rightarrow E_p(\omega) = E_0$$

$$S_{\text{dir}}(\omega) = -\frac{4\pi\omega |d|^4 \Gamma (1 + e^{-kT})}{m(\omega) (\Gamma^2 - (\omega - \omega_0)^2)} \quad \leftarrow \text{line shape} \quad E_0^2 E_{\text{pump}}^2 E_0^2$$



Two-dimensional lineshapes

$$E(t, T, \tau) = i\omega P(t, T, \tau) = i\omega S^{(R)}(t, T, \tau) e^{-i\omega(t-\tau)} \Theta(\tau) \\ + i\omega S^{(NR)}(t, T, \tau) e^{-i\omega(t-\tau)} \Theta(-\tau)$$

$$\tilde{E}(\omega_t, T, \omega_\tau) \approx \int_0^\infty dt \int_{-\infty}^\infty d\tau \left[S^{(R)}(t, T, \tau) e^{i\omega_t t - i\omega_\tau \tau} e^{-i\omega(t-\tau)} \right. \\ \left. + S^{(NR)}(t, T, \tau) e^{i\omega_t t - i\omega_\tau \tau} e^{-i\omega(t-\tau)} \right]$$

$$\approx \int_0^\infty dt \int_0^\infty d\tau e^{+i(\omega_t - \omega)t} e^{-i(\omega_\tau - \omega)\tau} e^{-\Gamma t} e^{-\Gamma \tau}$$

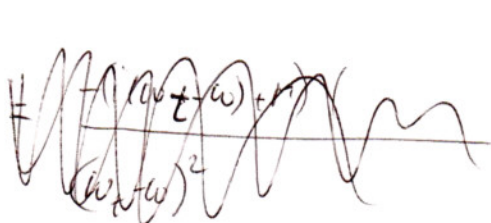


$$+ \int_0^\infty dt \int_{-\infty}^0 d\tau e^{i(\omega_t - \omega)t} e^{-i(\omega_\tau - \omega)\tau} e^{-\Gamma t} e^{-\Gamma|\tau|}$$

$$= \int_0^\infty dt \int_0^\infty d\tau \left(\frac{1}{i(\omega_t - \omega) - \Gamma} \left[e^{-\Gamma t} - e^{-\Gamma \tau} \right] \right) \left(\frac{1}{-i(\omega_\tau - \omega) - \Gamma} \left[e^{-\Gamma \tau} - e^{-\Gamma t} \right] \right)$$

$$+ \int_0^\infty dt \int_{-\infty}^0 d\tau' e^{i(\omega_t - \omega)t} e^{+i(\omega_\tau - \omega)\tau'} e^{-\Gamma t} e^{-\Gamma|\tau'|}$$

$$= \left(-\frac{1}{i(\omega_t - \omega) - \Gamma} \right) \left(\frac{1}{i(\omega_\tau - \omega) + \Gamma} \right) + \frac{+1}{i(\omega_t - \omega) - \Gamma} \frac{1}{i(\omega_\tau - \omega) - \Gamma}$$



↑
rephasing part

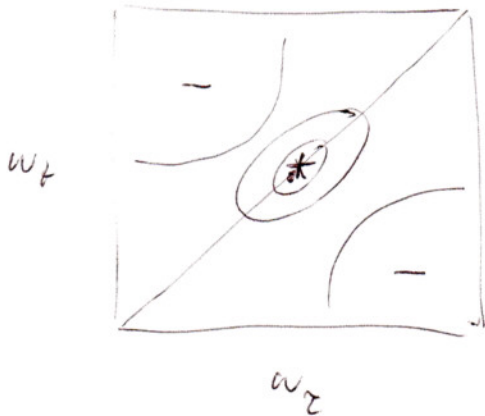
↑
non-rephasing part

Repeating part

$$(q+is)(q-is) = q^2 - s^2$$

$$\frac{-r - i(\omega_t - \omega)}{r^2 - (\omega_t - \omega)^2} \cdot \frac{r - i(\omega_r - \omega)}{r^2 - (\omega_r - \omega)^2} =$$

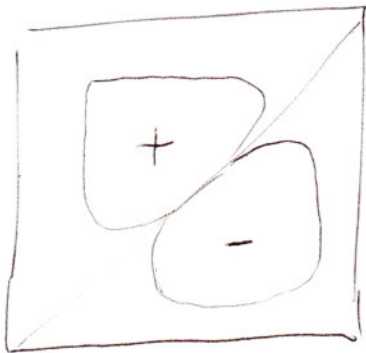
$$= \frac{r + i(\omega_t + \omega)}{r^2 - (\omega_t - \omega)^2} \cdot \frac{r - i(\omega_r - \omega)}{r^2 - (\omega_r - \omega)^2} = \frac{r^2 + (\omega_t - \omega)(\omega_r - \omega)}{(r^2 - (\omega_t - \omega)^2)(r^2 - (\omega_r - \omega)^2)}$$



$$+ i \frac{r[(\omega_t - \omega) - \omega_r + \omega]}{(r^2 - (\omega_t - \omega)^2)(r^2 - (\omega_r - \omega)^2)}$$

Real part

max at $\omega \equiv \omega_t$ and $\omega \equiv \omega_r$



Imaginary part

if $\omega_t \equiv \omega_r$ $I_m = 0$

if $\omega_t > \omega_r$ $I_m +$

Non-repeating part

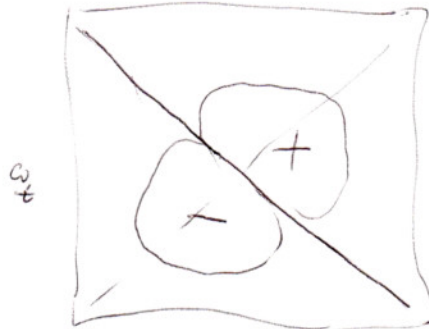
$$\frac{-r - i(\omega_t - \omega)}{r^2 - (\omega_t - \omega)^2} \cdot \frac{-r - i(\omega_r - \omega)}{r^2 - (\omega_r - \omega)^2} = \frac{r^2 + i r(\omega_t - \omega) + i r(\omega_r - \omega) - (\omega_t - \omega)(\omega_r - \omega)}{(r^2 - (\omega_t - \omega)^2)(r^2 - (\omega_r - \omega)^2)}$$

$$= \frac{r^2 - (\omega_r - \omega)(\omega_t - \omega)}{(r^2 - (\omega_t - \omega)^2)(r^2 - (\omega_r - \omega)^2)} + i \frac{r(\omega_t + \omega_r - 2\omega)}{(r^2 - (\omega_t - \omega)^2)(\dots)}$$

Real part



Imaginary part



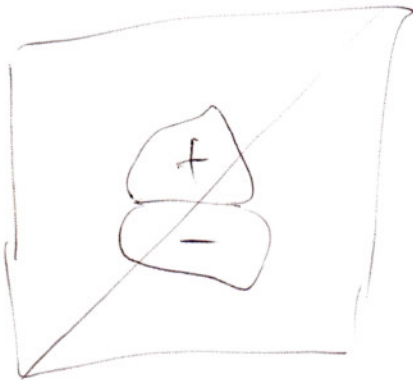
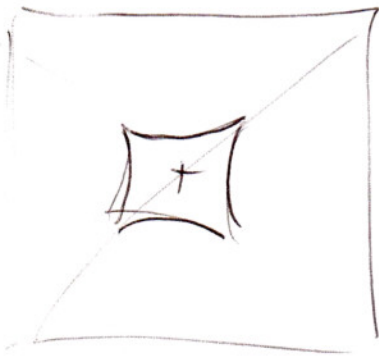
zero when $(\omega_t - \omega) = -(\omega_r - \omega)$

~~both $(\omega_t - \omega) > 0$~~

both $(\omega_t - \omega) > 0$
 $(\omega_r - \omega) > 0 \Rightarrow \text{Im} > 0$

both $(\omega_t - \omega) < 0$
 $(\omega_r - \omega) < 0 \Rightarrow \text{Im} < 0$

Complete spectra R-IR



The real part ~~is~~ taken at $(\omega_t - \omega) = 0$ gives

$$\frac{\Gamma^2}{(\Gamma^2 - (\omega_t - \omega)^2)(\Gamma^2 - (\omega_r - \omega)^2)} = \frac{1}{\Gamma^2 - (\omega_r - \omega)^2}$$

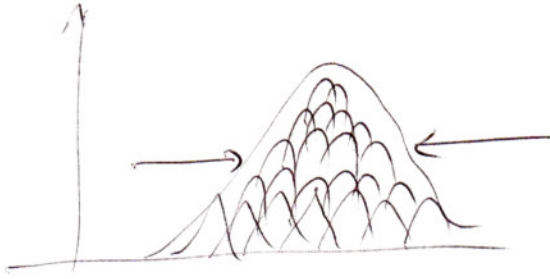
homogeneous line shape.

Taking anti-diagonal width just a little bit different, but approximately it gives the same result.

This remains true even if the system is inhomogeneously broadened.

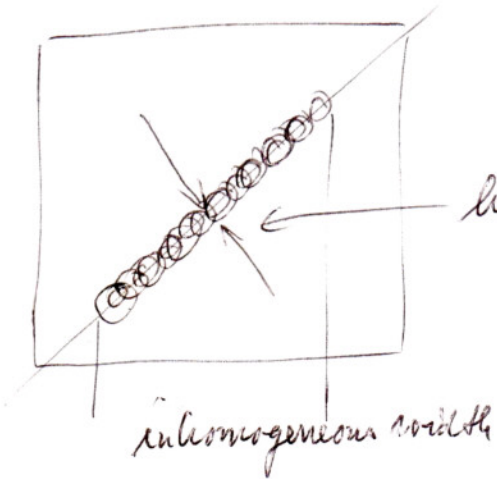
Absorption spectrum and inhomogeneous broadening

Abs



inhomogeneous width is \gg
than the homogeneous

2D and IB



homogeneous width is still visible!